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Fractional Calculus: Methods for Applications

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Topics to be discussed

- Interpretations of fractional-order operators
- Numerical fractional differentiation and numerical solution of FDEs
- Matrix approach: from CO to VO to DO
- Building fractional-order models: fitting using the Mittag-Leffler function
- "Least circles" method
- Related Matlab toolboxes
- Applications

















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Most used definitions of fractional differentiation

Caputo, 1967:

For $f \in C^{(n)}[a,b]$, $f^{(k)}(a) = 0$ (k = 0, ..., n - 1)R-L, C, M-R and G-L definitions

are equivalent

Miller-Ross, 1990s

(Dzhrbashyan, 1960s)

Grünwald–Letnikov, 1860s (Liouville, 1830s)

$$D^{\alpha}f(t) = \lim_{h \to 0} h^{-\alpha} \sum_{k=0}^{\left[\frac{t-\alpha}{h}\right]} (-1)^k \binom{\alpha}{n} f(t-kh)$$



$$\begin{array}{l} \textbf{Geometric interpretation}\\ \textbf{of fractional integration:}\\ \textbf{shadows on the walls} \end{array}$$
$${}_0I_t^{\alpha}f(t) = \frac{1}{\Gamma(\alpha)}\int_0^t f(\tau)(t-\tau)^{\alpha-1}d\tau, \quad t \ge 0, \\\\ {}_0I_t^{\alpha}f(t) = \int_0^t f(\tau) \, dg_t(\tau), \\\\ g_t(\tau) = \frac{1}{\Gamma(\alpha+1)} \{t^{\alpha} - (t-\tau)^{\alpha}\}. \end{array}$$
For $t_1 = kt, \ \tau_1 = k\tau \ (k > 0)$ we have:
 $g_{t_1}(\tau_1) = g_{kt}(k\tau) = k^{\alpha}g_t(\tau). \end{array}$









Physical interpretation of Stieltjes integral (1)

Imagine a car equipped with two devices for measurements: the speedometer recording the velocity $v(\tau)$, and the clock which should show the time τ . The clock, however, shows the time incorrectly.

Suppose that the relationship between the wrong time τ , shown by the clock and considered by the driver A as the correct time, and the true time T, on the other, is described by the function $T = g(\tau)$.

Physical interpretation of Stieltjes integral (2)

Driver A's computations:

$$S_A(t) = \int_0^t v(\tau) d\tau$$
.

Well-informed observer O's computations:

$$S_O(t) = \int_0^t v(\tau) dg(\tau).$$

This example shows that the Stieltjes integral can be interpreted as the real distance passed by a moving object, for which we have recorded correct values of speed and incorrect values of time; the relationship between the wrongly recorded time τ and the correct time T is given by a known function $T = g(\tau)$.

Physical interpretation of the Riemann-Liouville integral: "shadows of the past"

$$\begin{split} S_O(t) &= \int\limits_0^t v(\tau) \, dg_t(\tau) = \ _0 I_t^\alpha v(t), \\ g_t(\tau) &= \frac{1}{\Gamma(\alpha+1)} \{ t^\alpha - (t-\tau)^\alpha \}. \end{split}$$

The left-sided Riemann–Liouville fractional integral of the individual speed $v(\tau)$ of a moving object, for which the relationship between its individual time τ and the cosmic time T at each individual time instance t is given by the known function $T = g_t(\tau)$, represents the real distance $S_O(t)$ passed by that object.

Interpretations of the Volterra integrals

$$K * f(t) = \int_{0}^{t} f(\tau)k(t-\tau)d\tau$$

Assuming that k(t) = K'(t), this integral takes the form:

$$K * f(t) = \int_{0}^{t} f(\tau) dq_t(\tau),$$

 $q_t(\tau) = K(t) - K(t - \tau).$ The geometric and physical interpretations of the Volterra convolution integral are then similar to the suggested interpretations for fractional integrals.

Physical interpretation of initial conditions for fractional differential equations with the Riemann-Liouville fractional derivatives

$$\begin{split} \int_{0}^{\infty} e^{-st} \,_{0} D_{t}^{\alpha} f(t) \, dt &= s^{\alpha} F(s) - \sum_{k=0}^{n-1} s^{k} \, \left[{}_{0} D_{t}^{\alpha-k-1} f(t) \right]_{t=0}, \\ & (n-1 < \alpha \le n). \end{split}$$

$$t^{\alpha k+\beta-1} E_{\alpha,\beta}^{(k)}(\pm z t^{\alpha}) \quad \leftrightarrows \quad \frac{k! \, p^{\alpha-\beta}}{(p^{\alpha} \mp a)^{k+1}}$$

The Laplace transform method for FDEs

EXAMPLE I.

$$_{0}D_{t}^{1/2}f(t) + af(t) = 0, \quad (t > 0); \qquad \left[{}_{0}D_{t}^{-1/2}f(t) \right]_{t=0} = C$$

Solution: applying the Laplace transform we obtain

$$F(s) = \frac{C}{s^{1/2} + a}, \qquad C = \left[{}_{0}D_{t}^{-1/2}f(t) \right]_{t=0}$$

and the inverse Laplace transform gives the solution:

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$$f(t) = Ct^{-1/2} E_{\frac{1}{2},\frac{1}{2}}(-a\sqrt{t}).$$

If a=1, then

$$f(t) = C(\frac{1}{\sqrt{\pi t}} - e^t \operatorname{erfc}(\sqrt{t}))$$

The Laplace transform method for FDEs EXAMPLE 2.

 $\begin{array}{c} 0 D_t^{\alpha} y(t) - \lambda y(t) = h(t), \quad (t > 0); \\ \left[0 D_t^{\alpha - k} y(t) \right]_{t=0} = b_k, \quad (k = 1, 2, \dots, n) \end{array}$

$$n-1 < \alpha < n.$$

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Solution: applying the Laplace transform:

$$s^{\alpha}Y(s) - \lambda Y(s) = H(s) + \sum_{k=1}^{n} b_k s^{k-1},$$
$$Y(s) = \frac{H(s)}{s^{\alpha} - \lambda} + \sum_{k=1}^{n} b_k \frac{s^{k-1}}{s^{\alpha} - \lambda}$$

and the inverse transform gives the solution:

$$y(t) = \sum_{k=1}^{n} b_k t^{\alpha-k} E_{\alpha,\alpha-k+1}(\lambda t^{\alpha}) + \int_0^t (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^{\alpha})h(\tau)d\tau$$

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We deal with the Riemann-Liouville derivatives $(n - 1 \le \alpha < n)$:

$${}_0D_t^{\alpha}f(t) = \frac{1}{\Gamma(n-\alpha)} \left(\frac{d}{dt}\right)^n \int_0^t \frac{f(\tau) d\tau}{(t-\tau)^{\alpha-n+1}}$$

+

Fractional differential equations in terms of RL derivatives require initial conditions expressed in terms of initial values of fractional derivatives of the unknown function.

A typical initial value problem $(n - 1 < \alpha < n)$:

$$\begin{split} &_{0}D_{t}^{\alpha}f(t)+af(t)=h(t); \qquad (t>0)\\ & \left[{}_{0}D_{t}^{\alpha-k}f(t)\right]_{t\rightarrow0}=b_{k}, \qquad (k=1,2,\ldots,n). \end{split}$$

What about initial conditions?

$$ig[{}_0D_t^{lpha-k}y(t)ig]_{t=0}=b_k,\qquad (k=1,\ldots,[lpha]+1)$$

K. Diethelm, N. J. Ford, A. D. Freed, and Yu. Luchko (January 2005):

"A typical feature of differential equations (both classical and fractional) is the need to specify additional conditions in order to produce a unique solution. For the case of Caputo FDEs, these additional conditions are just the static initial conditions ..., which are akin to those of classical ODEs, and are therefore familiar to us. In contrast, for Riemann-Louville FDEs, these additional conditions constitute certain fractional derivatives (and/or integrals) of the unknown solution at the initial point $x=0\ldots$, which are functions of x. These initial conditions are not physical; furthermore, it is not clear how such quantities are to be measured from experiment, say, so that they can be appropriately assigned in an analysis."

N. Heymans and I. Podlubny:

"Physical interpretation of initial conditions for fractional differential equations with the Riemann-Liouville fractional derivatives", *Rheologica Acta*, vol. 45, no. 5, June 2006, pp. 765–772.

Spring-pot model: Creep

A stress step σ_0 is applied at initial time t=0. The change of $\epsilon(t)$ is described by the FDE

$$_0 D_t^{\alpha} \epsilon(t) = \frac{\sigma_0}{K}$$

An initial condition involving $_0D_t^{\alpha-1}\epsilon(t)$ is required. It can be found by taking the first-order integral of the constitutive equation and letting $t\to 0$

 $\left[{}_0D_t^{\alpha-1}\epsilon(t)\right]_{t\to 0} = \left[{}_0D_t^{-1}(\sigma_0/K)\right]_{t\to 0}.$

In the considered case stress is finite at all times, therefore the required IC is

 $\left[{}_{0}D_{t}^{\alpha-1}\epsilon(t)\right]_{t\to0}=0.$

Spring-pot model

Spring-pot is a linear viscoelastic element whose behaviour is intermediate between that of elastic element (spring) and a viscous element (dashpot). The term "spring-pot" was introduced by Koeller (1984), although the concept of an element with intermediate properties had been introduced some time earlier (G. W. Scott Blair, 1930s-40s). The constitutive equation of a spring-pot is:

$$\sigma(t) = K_0 D_t^{\alpha} \epsilon(t)$$
 or $\epsilon(t) = \frac{1}{K} {}_0 D_t^{-\alpha} \sigma(t)$

Spring-pot model: Impulse response

An impulse of stress defined as $B\delta(t)$ applied to the spring-pot at time t = 0. After that, the stress remains zero. The strain $\epsilon(t)$ for t > 0 is the solution of FDE

 ${}_{0}D_{t}^{\alpha}\epsilon(t)=0.$

An initial condition involving $\left[{}_0 D_t^{lpha-1} \epsilon(t)
ight]_{t
ightarrow 0}$ is required.

This can be found through integration of the constitutive equation,

 $\left[{}_0D_t^{\alpha-1}\epsilon(t)\right]_{t\to0}=\left[{}_0D_t^{-1}\sigma(t)/K\right]_{t\to0},$ which gives the following initial condition:

as

$$\left[{}_{0}D_{t}^{\alpha-1}\epsilon(t) \right]_{t \to 0} = B/K.$$

The key: look for inseparable twins

In a general case, when we consider some FDE for, say, U(t), we have to consider also some function V(t), for which some dual relation exists between U(t) and V(t). For example: stress $\sigma(t)$ and strain $\epsilon(t)$ in viscoelasticity; charge q(t) and voltage v(t) in electrical circuits; temperature difference T(t) and the heat flux q(t) in heat conduction; etc. Functions U(t) and V(t) are normally related by some basic physical law for the particular field of science.

In each scientific field there are such pairs of functions like the aforementioned, which are as *inseparable as Siamese twins*: the left-hand side of the initial condition involves one of them, whereas the evaluation of the right-hand side is related to the other.